

# Chirality in soft matter: from out-of-equilibrium physics to non-linear optics

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Fakulteta za *matematiko in fiziko*



**arrs**

JAVNA AGENCIJA ZA RAZISKOVALNO DEJAVNOST  
REPUBLIKE SLOVENIJE

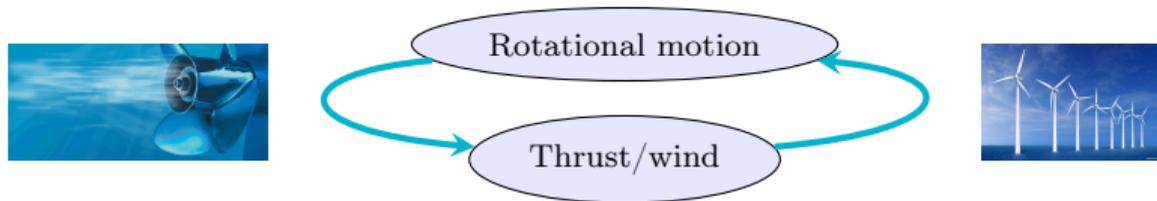


# Outline

- 1 Introduction
- 2 Lehmann effect: an out-of-equilibrium effect in chiral liquid crystal droplets
- 3 Interlude: light propagation in anisotropic media
- 4 Role of chirality in the non-linear response of a confined cholesteric

# Chirality in everyday life

- Chiral object: distinguishable from its mirror image.
- A common example: propeller.



- Without chirality, this conversion is not possible.

## Chirality in soft matter: the cholesteric phase

- Nematic liquid crystal: no positional order, mean molecular orientation  $\mathbf{n}$

## Chirality in soft matter: the cholesteric phase

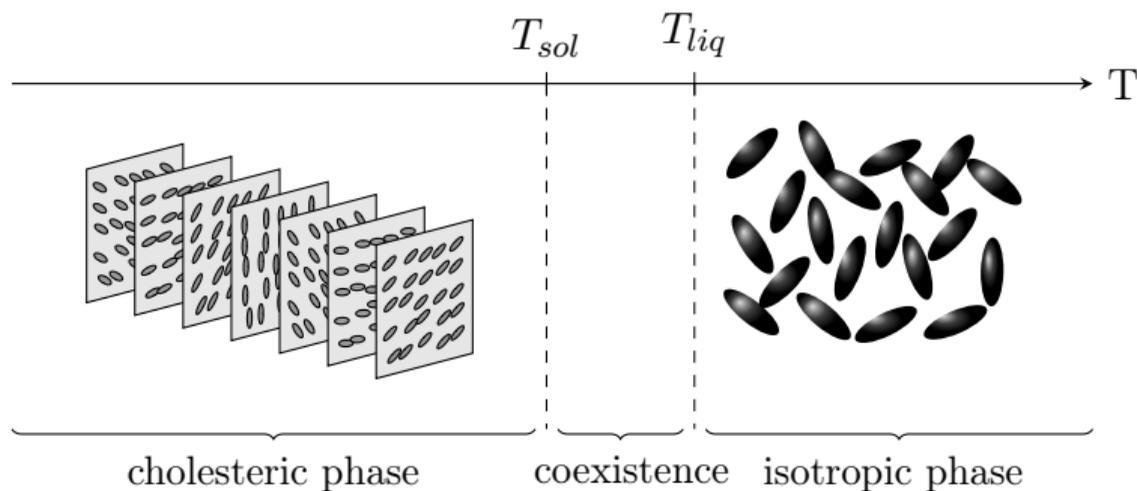
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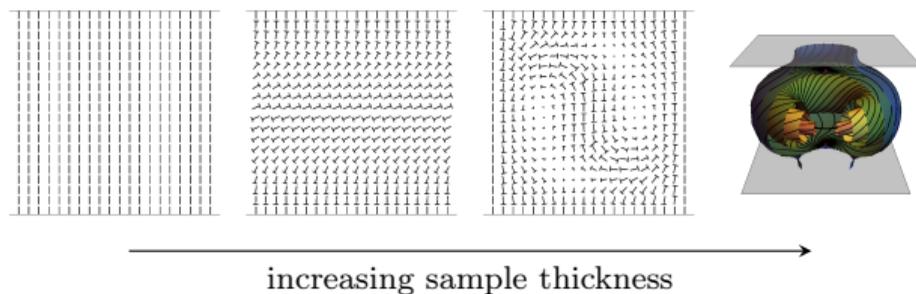
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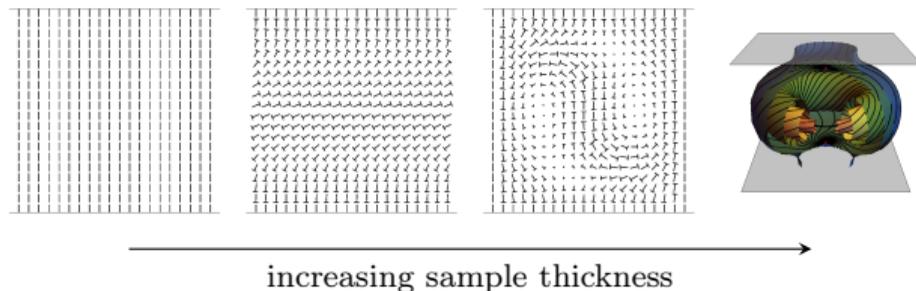
## Confining cholesterics between two plates

- Surface constraint: molecules must be normal to the confining surface

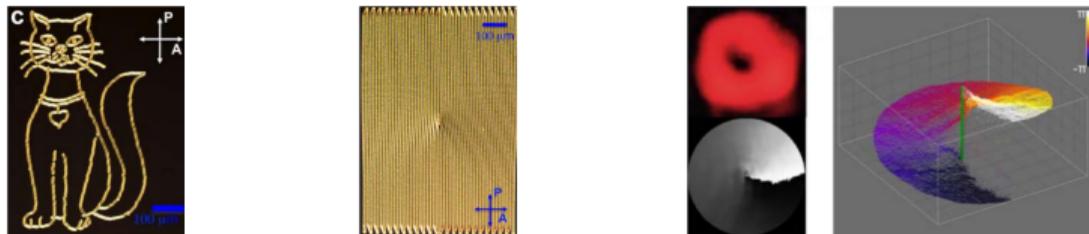


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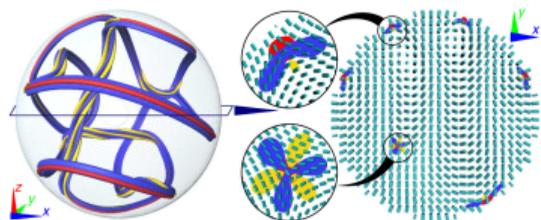
- Arbitrary shapes can be written!



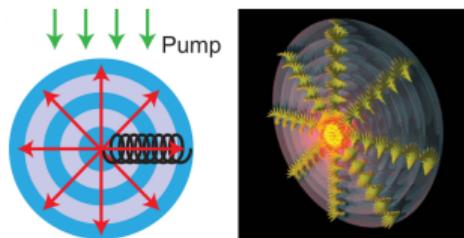
P. J. Ackerman et al. *Scientific Reports*, 2, 2012

# Confining cholesterics inside droplets

Topological zoo of free standing knots



Lasing in a cholesteric droplet: an omnidirectional microscopic coherent light source

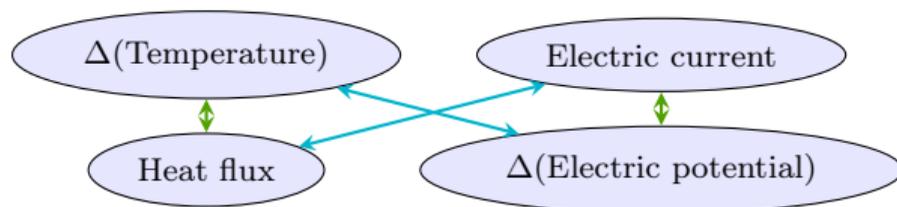


D. Seč, S. Čopar, and S. Žumer. *Nature Communications*, 5:3057, 2014

M. Humar. *Liquid Crystals*, 43:1937–1950, 2016

## Other aspects of chirality in soft matter

Cross-coupling effects in out-of-equilibrium systems:



Applications:

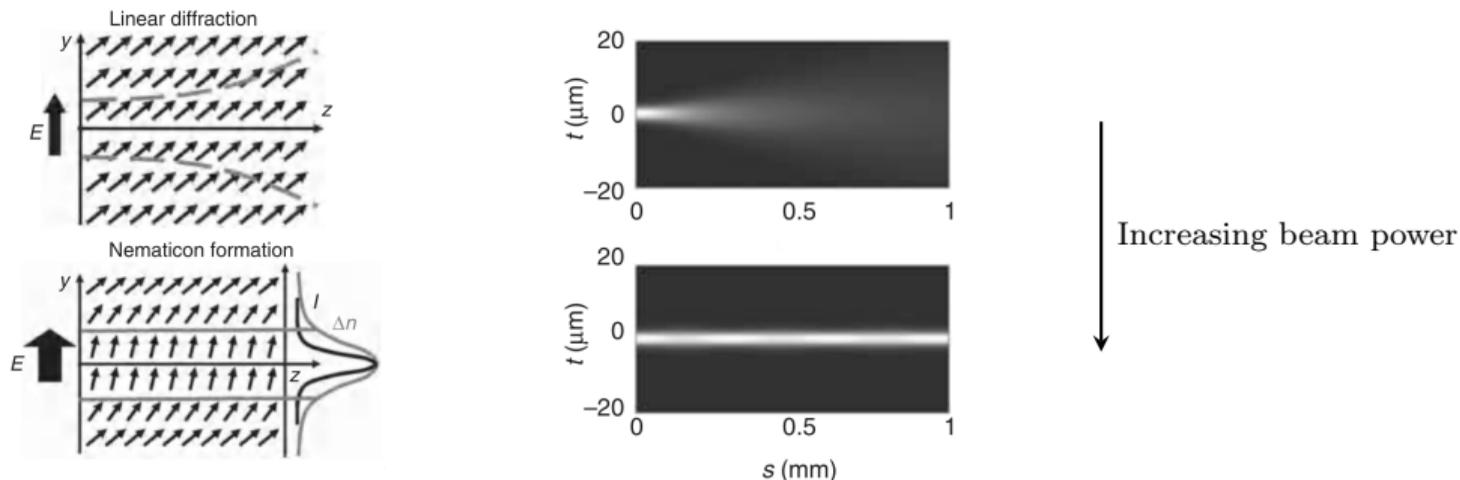


## Problematic

Role of chirality in confined liquid-crystal systems submitted to a temperature gradient?

## Other aspects of chirality in soft matter

Non-linear optical response of liquid crystal systems:



G. Assanto. *Nematicons*. John Wiley & Sons, 2013

## Problematic

Role of chirality in the non-linear optical response of a confined cholesteric?

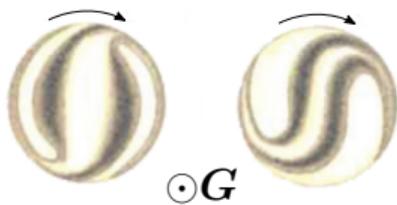
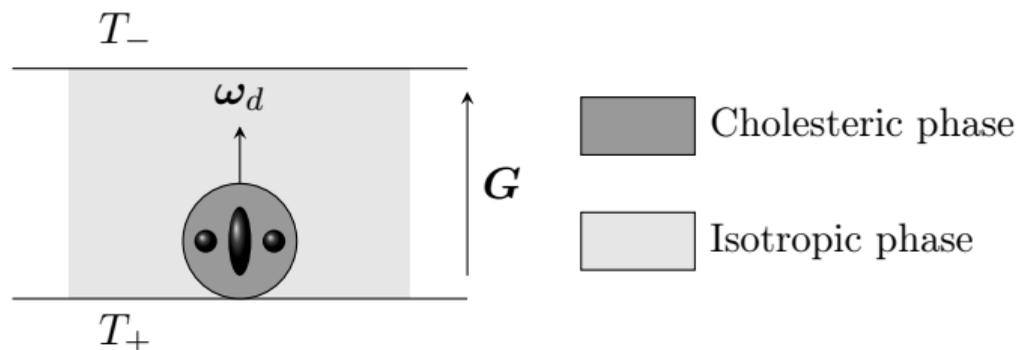
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## First observations by Lehmann



Lehmann, 1900:

- coexistence of cholesteric droplets with the isotropic fluid
- rotation of the droplets internal texture when heated from below

## What is causing the rotation of the Lehmann droplets

Rotation because of the microscopic or macroscopic chirality?

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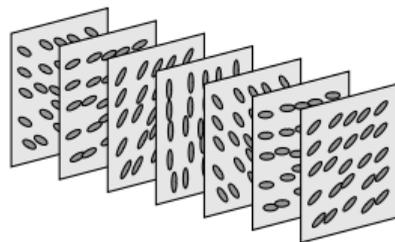
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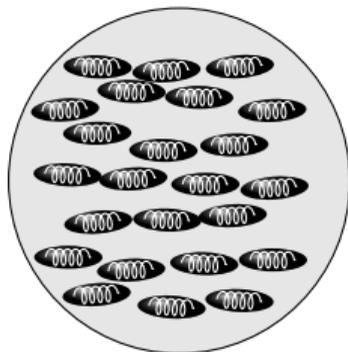


- macroscopic chirality  $\Leftrightarrow$  twisted texture (helix in at least one direction)



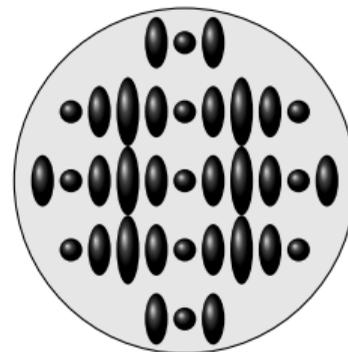
# What is causing the rotation of the Lehmann droplets

Possible tests:



{ chiral molecules  $\leftrightarrow$  cholesteric  
no macroscopic twist (compensated)

Thermal gradient  $\Rightarrow$  no rotation



{ no chiral molecules  $\leftrightarrow$  nematic  
macroscopic twist

Thermal gradient  $\Rightarrow$  rotation?

## Question

Can we observe the Lehmann effect in droplets of a **nematic achiral phase** with a **chiral director field**?

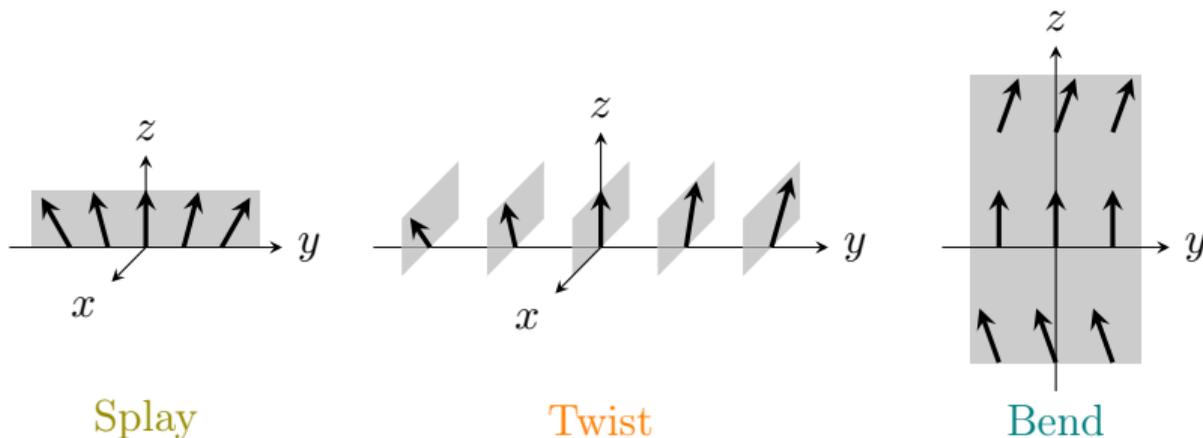
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## Elastic deformations in a nematic phase

Frank-Oseen elastic energy:

$$F[\mathbf{n}] = \int_V \frac{dV}{2} (K_1 [\nabla \cdot \mathbf{n}]^2 + K_2 [\mathbf{n} \cdot \nabla \times \mathbf{n}]^2 + K_3 [\mathbf{n} \times \nabla \times \mathbf{n}]^2)$$



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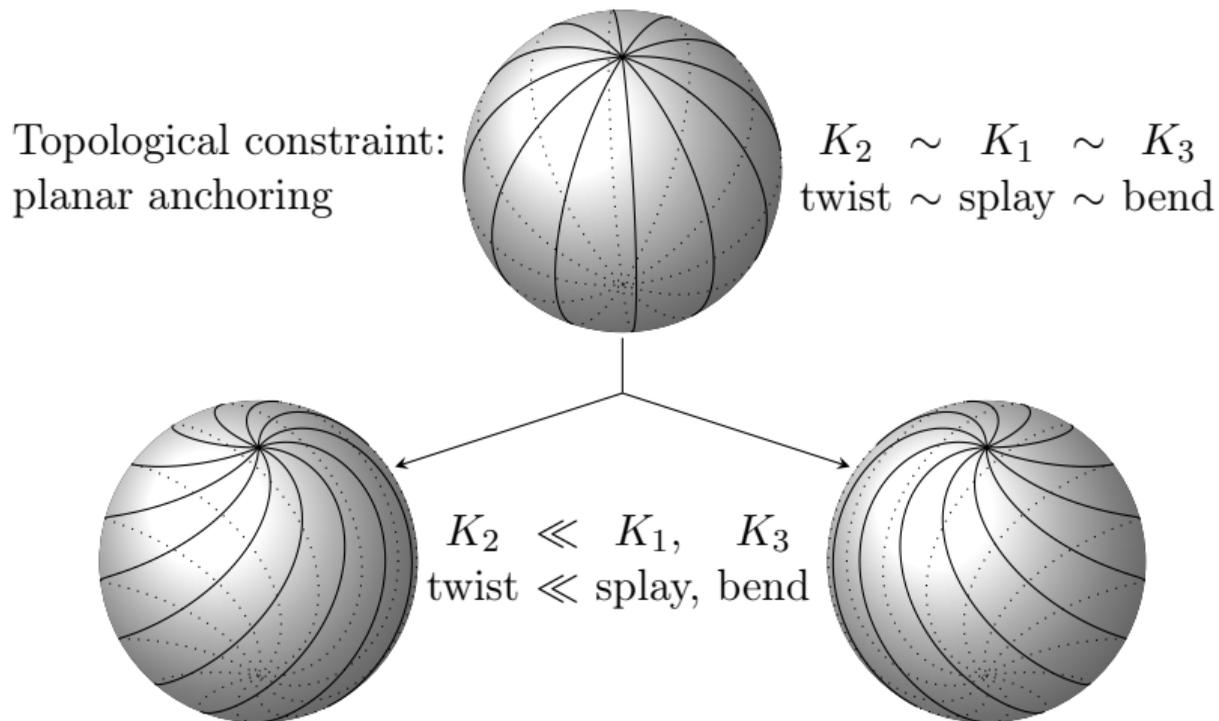
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- action of a **topological constraint** on the LC domain surface:
  - ★  $F[\mathbf{n}] \rightarrow F[\mathbf{n}] + \int_S dS \gamma(\mathbf{n})$ , with  $\gamma$  the anchoring energy

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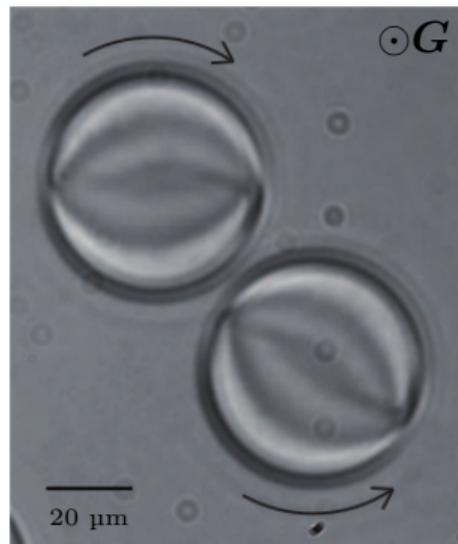
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## Stability of bipolar configuration



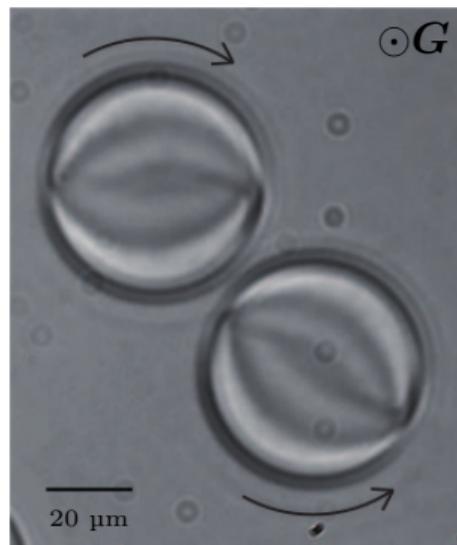
R. D. Williams. *Journal of physics A: mathematical and general*, 19:3211, 1986

## Rotation of twisted bipolar droplets



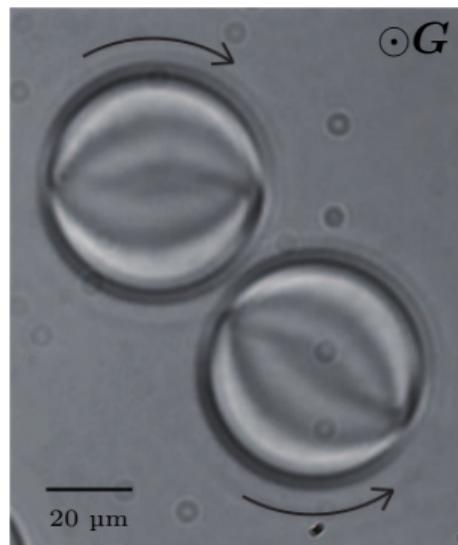
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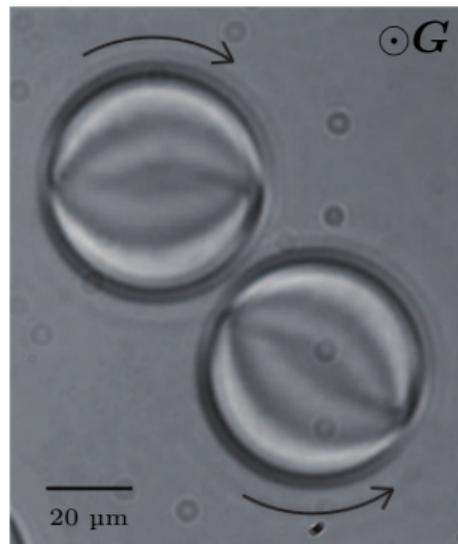
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Rotation only due to the twist of the director field

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- Lehmann effect in an achiral phase with a twisted director field:

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J. Ignés-Mullol, G. Poy, and P. Oswald. *Physical Review Letters*, 117:057801, 2016

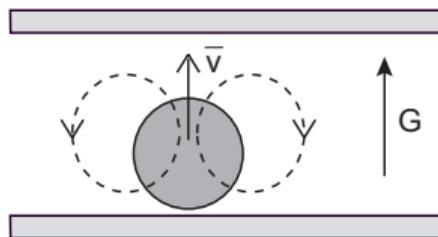
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- What is the "right" mechanism behind the Lehmann effect?



Melting-growth model: a gradient of impurity drives the molecules upward inside the droplet while the droplet interface stays fixed

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  - ★ Other methods (in-house implementation)

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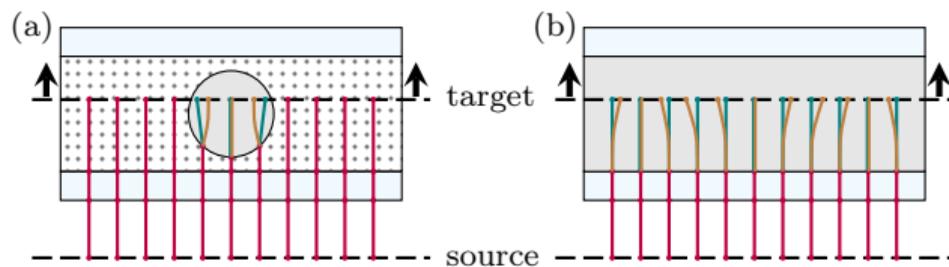
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Need for advanced light propagation code, if possible open-source

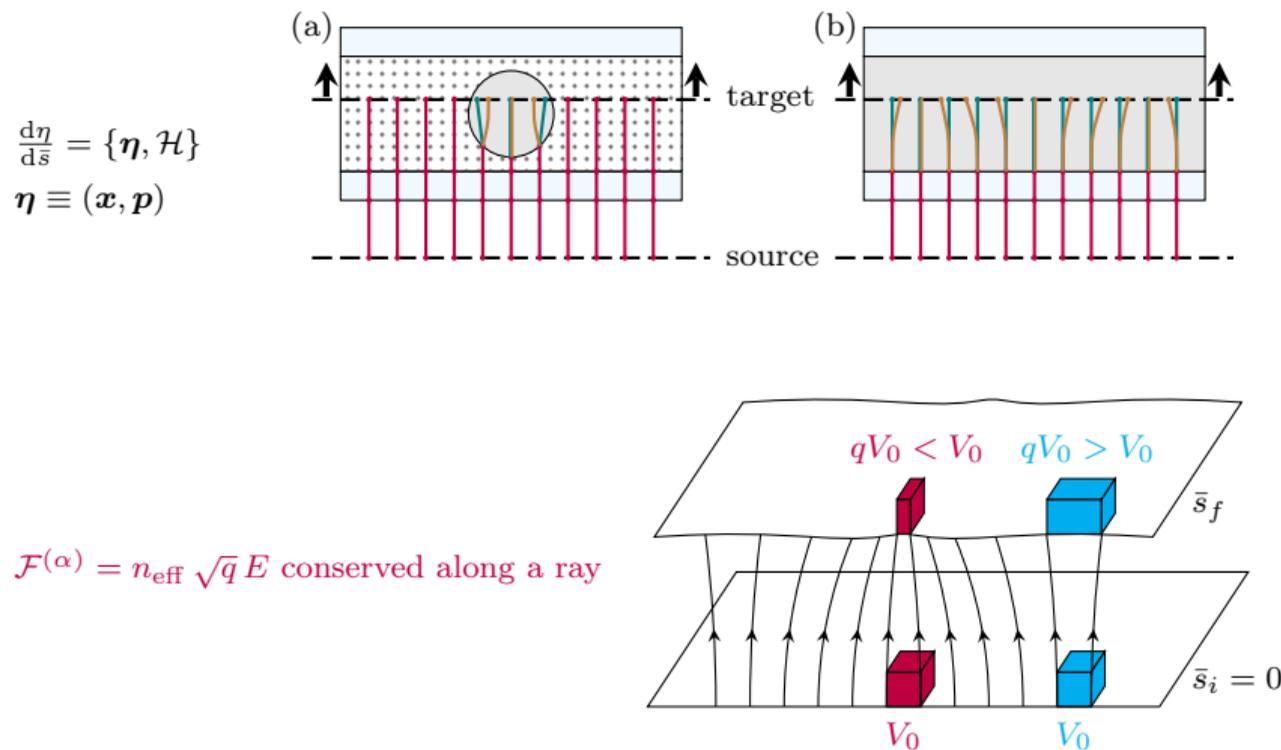
## First approach: Hamiltonian ray-tracing and energy transport

$$\frac{d\eta}{ds} = \{\eta, \mathcal{H}\}$$

$$\eta \equiv (\mathbf{x}, \mathbf{p})$$



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G. Poy and S. Žumer. *Soft Matter*, 15:3659–3670, 2019

## Second approach: physics-based splitting of the wave equation

- Wave-equation in anisotropic media:  $[\partial_k \partial_k \delta_{ij} - \partial_i \partial_j + k_0^2 \epsilon_{ij}] E_j = 0$

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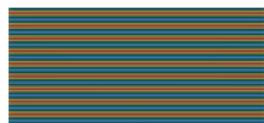
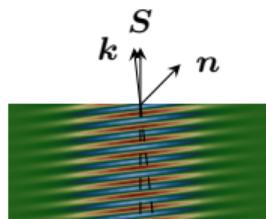
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$$i\partial_z \mathbf{E}_\perp = -\mathcal{P} \mathbf{E}_\perp$$

- What's inside  $\mathcal{P}$ ?

Phase op.  $\mathbf{K} \sim k_0^2 \epsilon$ Walkoff op.  $\mathbf{W} \sim (\epsilon \mathbf{u}_z) \otimes \nabla_\perp$ Diffraction op.  $\mathbf{D} \sim \Delta_\perp$

## Nemaktis: an open-source package for polarised microscopy

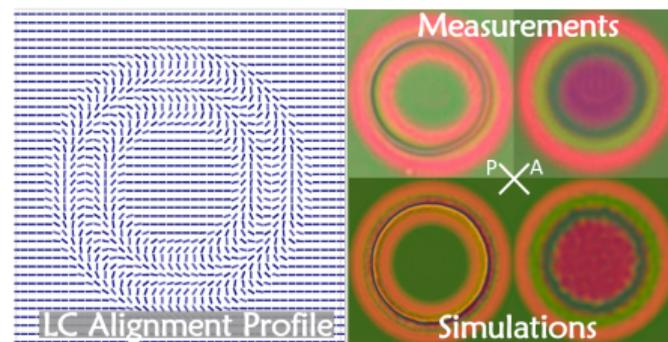
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- Closed-source BPM code for advanced uses: wide-angle beam deflection, non-linear optics, etc.



B. Berteloot, I. Nys, G. Poy, J. Beekman, and K. Neyts. *Soft Matter*, 16:4999, 2020

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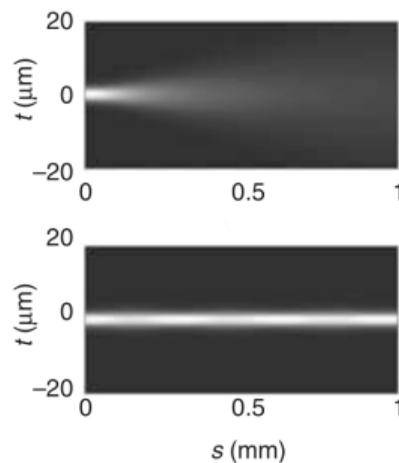
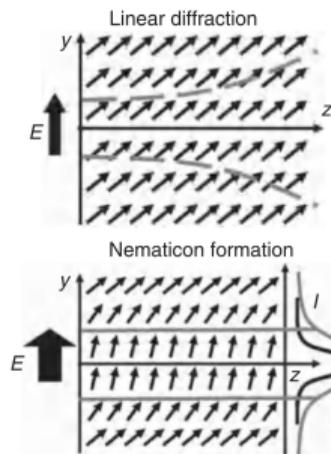
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## Motivations

## Spatial light solitons in liquid crystals: nematicons

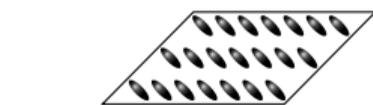


Increasing beam power

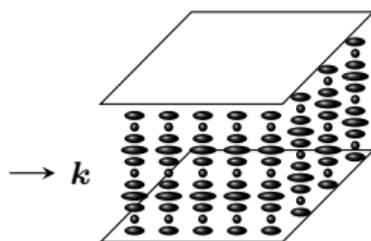
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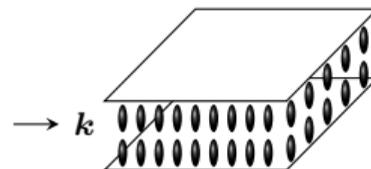
Studied systems in the past 20 years:



thick samples  
with planar  $n$



thick samples with  
cholesteric helix

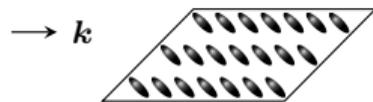
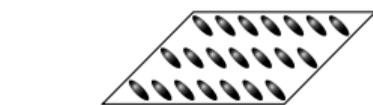


thin samples with  
homeotropic  $n$

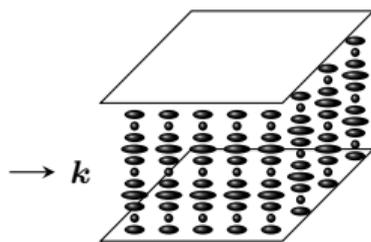
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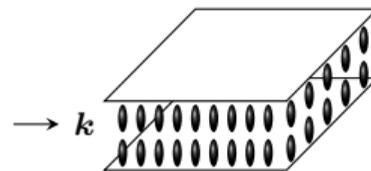
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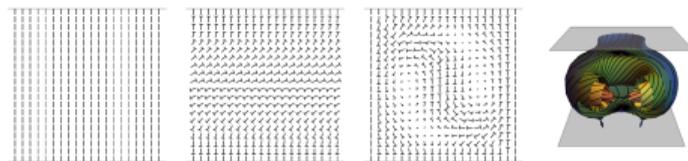
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# Motivations

What makes frustrated cholesteric (FCLC) an interesting system:

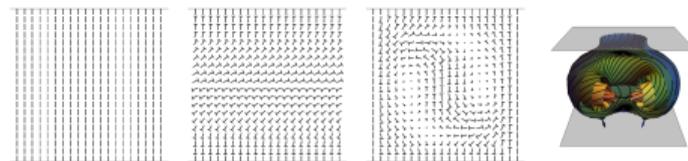
- Metastability for carefully chosen values of  $d/P$



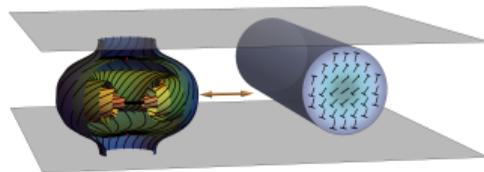
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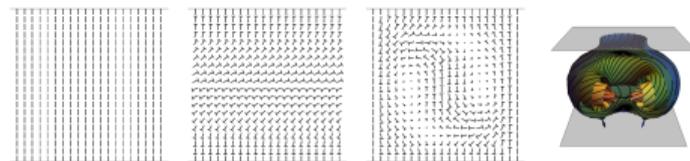
- Rich possibilities of interaction between light beams and topological solitons.



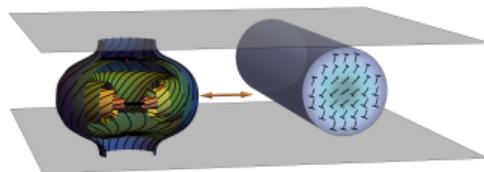
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### Problematic

Can we generate light solitons in frustrated cholesteric?

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## Orientational elasticity and non-linear interactions

Free energy of the liquid crystal phase:

$$F[\mathbf{n}, \mathbf{E}] = \int_V dV \left[ f_F(\mathbf{n}, \nabla \mathbf{n}) - \frac{\epsilon_0 \epsilon_a |\mathbf{n} \cdot \mathbf{E}|^2}{4} \right]$$

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Non-linear iterative scheme:

- $\mathbf{E}_{k+1}$ : BPM solution with  $\boldsymbol{\epsilon} = \epsilon_{\perp} \mathbf{I} + \epsilon_a \mathbf{n}_k \otimes \mathbf{n}_k$
- $\mathbf{n}_{k+1} = \mathbf{n}_k + \mu \frac{\delta F}{\delta \mathbf{n}} [\mathbf{n}_k, \mathbf{E}_{k+1}]$

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Free energy of the liquid crystal phase:

$$F[\mathbf{n}, \mathbf{E}] = \int_V dV \left[ f_F(\mathbf{n}, \nabla \mathbf{n}) - \frac{\epsilon_0 \epsilon_a |\mathbf{n} \cdot \mathbf{E}|^2}{4} \right]$$

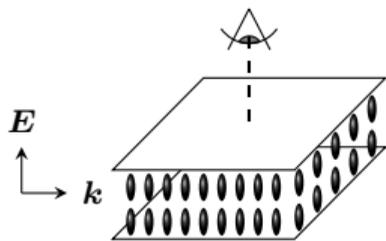
Non-linear iterative scheme:

- $\mathbf{E}_{k+1}$ : BPM solution with  $\boldsymbol{\epsilon} = \epsilon_{\perp} \mathbf{I} + \epsilon_a \mathbf{n}_k \otimes \mathbf{n}_k$
- $\mathbf{n}_{k+1} = \mathbf{n}_k + \mu \frac{\delta F}{\delta \mathbf{n}} [\mathbf{n}_k, \mathbf{E}_{k+1}]$

Typical running time for a mesh of  $3 \times 10^6$  points: **4 s / step**

(Full resolution of Maxwell equations for the same mesh:  $\sim 1$  h)

## Top-view observations



Top view of the thickness-averaged laser intensity (simulation):

Linear optical regime



Non-linear optical regime

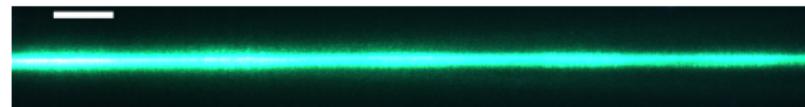


Top view of the scattered laser light (experiments):

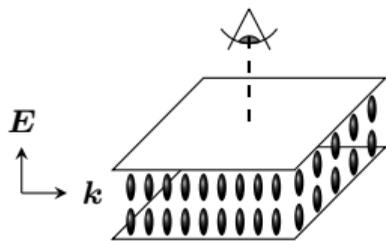
Linear optical regime



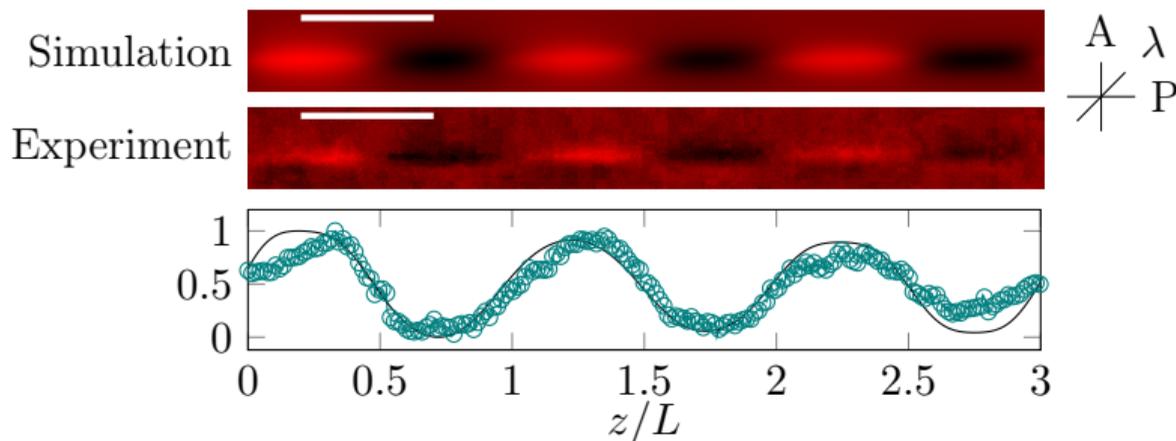
Non-linear optical regime



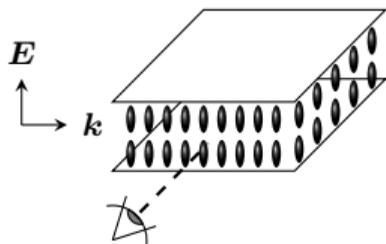
## Top-view observations



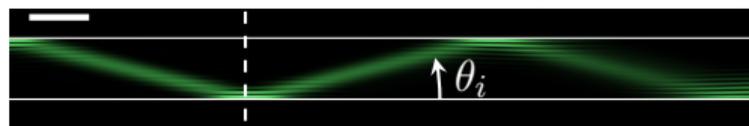
Top view polarised optical micrograph:



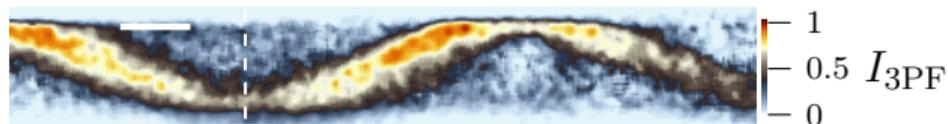
# Why is there a periodic molecular reorientation?



Side slice of beam intensity (simulation):

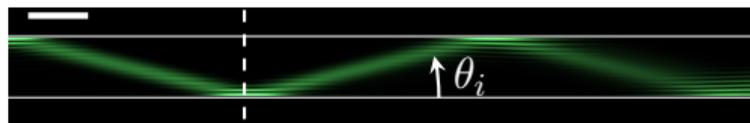


Side slice of 3PF signal (experiment):

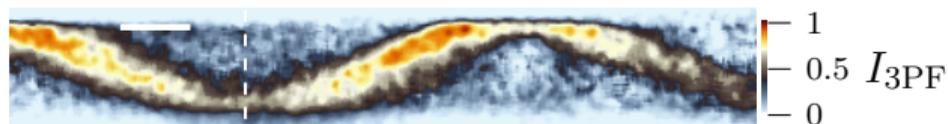


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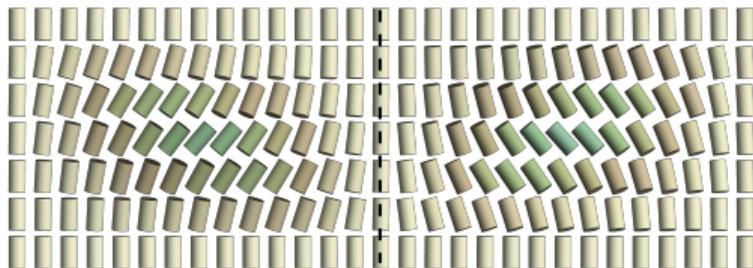
Side slice of beam intensity (simulation):



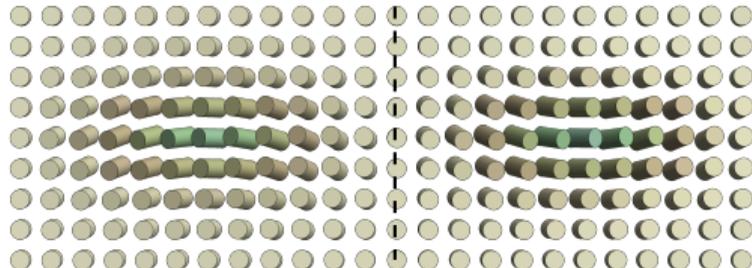
Side slice of 3PF signal (experiment):



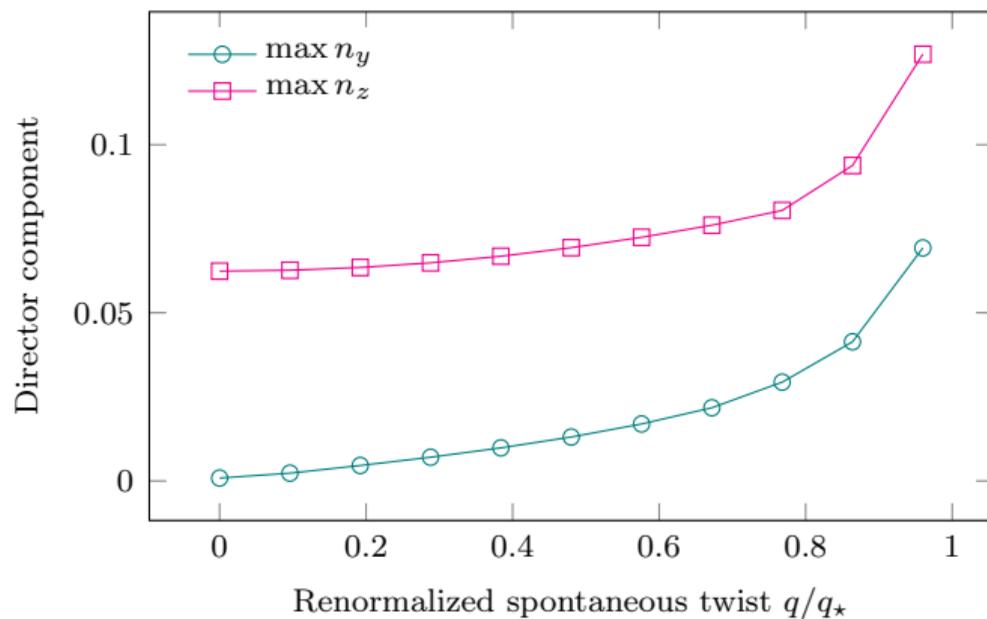
Side slice of director field



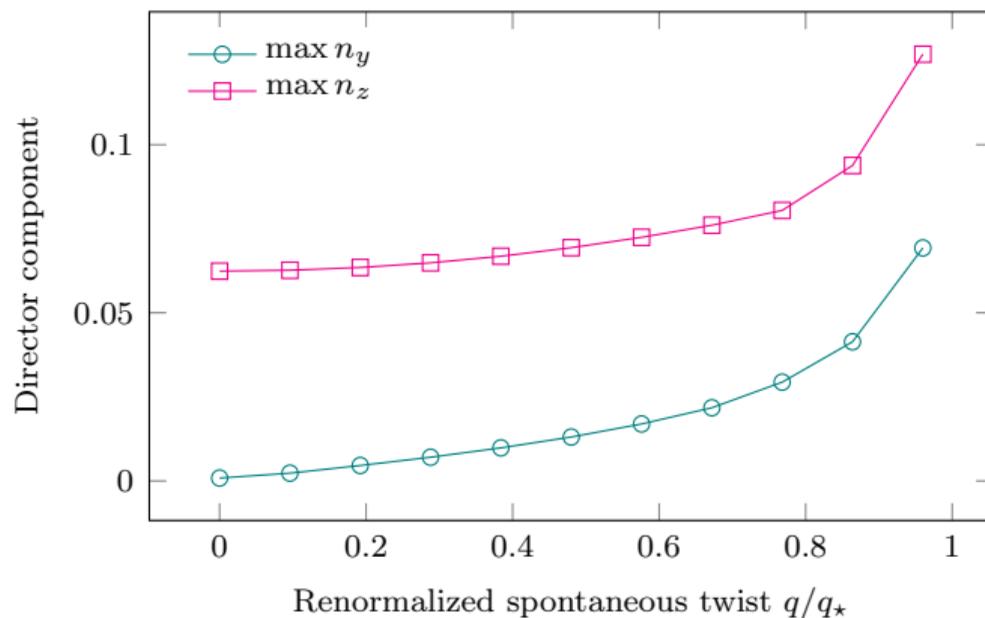
Mid-sample slice of director field



## Chirality-enhanced non-linear optical response



## Chirality-enhanced non-linear optical response



⇒ Potential for low-power non-linear optical photonics devices (e.g. active lenses)

# Outline

- 1 Introduction
- 2 Lehmann effect: an out-of-equilibrium effect in chiral liquid crystal droplets
- 3 Interlude: light propagation in anisotropic media
- 4 Role of chirality in the non-linear response of a confined cholesteric
  - Motivations
  - Light solitons in frustrated cholesteric
  - **Summary**

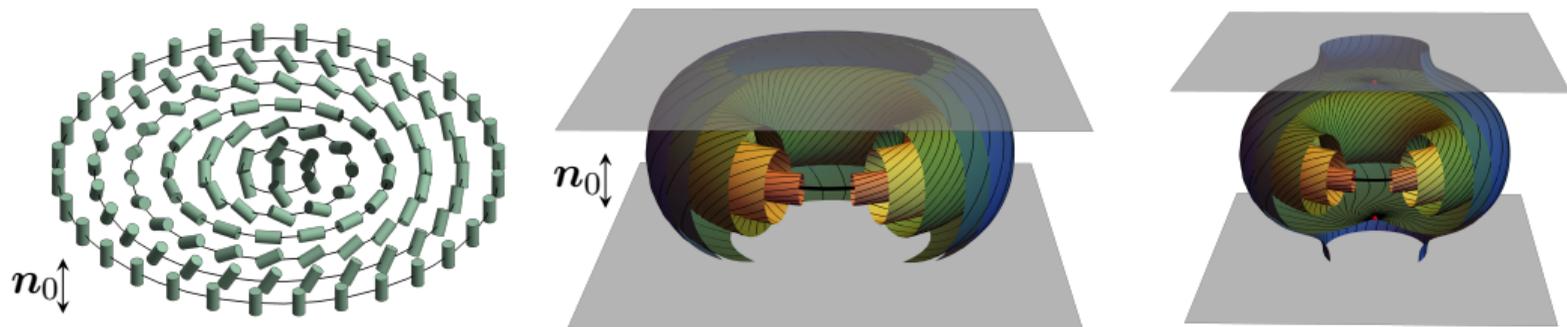
## Summary

- It is possible to generate solitons in confined cholesteric system, with:
  - ★ "bouncing" beam between the sample plates
  - ★ periodic reorientation along the beam axis
  - ★ chirality-enhanced Kerr response

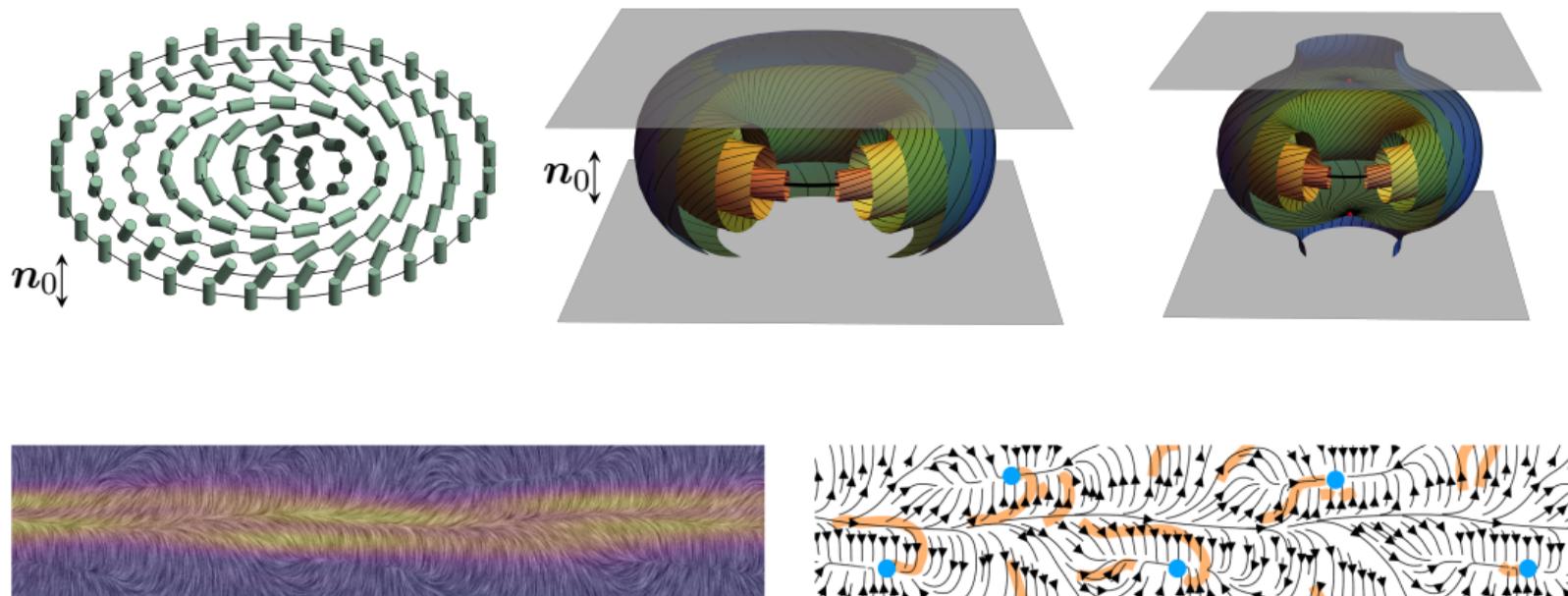
## Summary

- It is possible to generate solitons in confined cholesteric system, with:
  - ★ "bouncing" beam between the sample plates
  - ★ periodic reorientation along the beam axis
  - ★ chirality-enhanced Kerr response
- To be explored:
  - ★ Superposition of normal and transverse polarisations (spin-orbit solitons)
  - ★ Interaction with topological solitons (topological optomechanics)

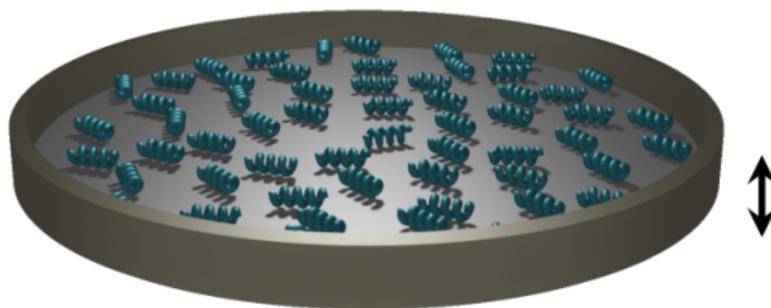
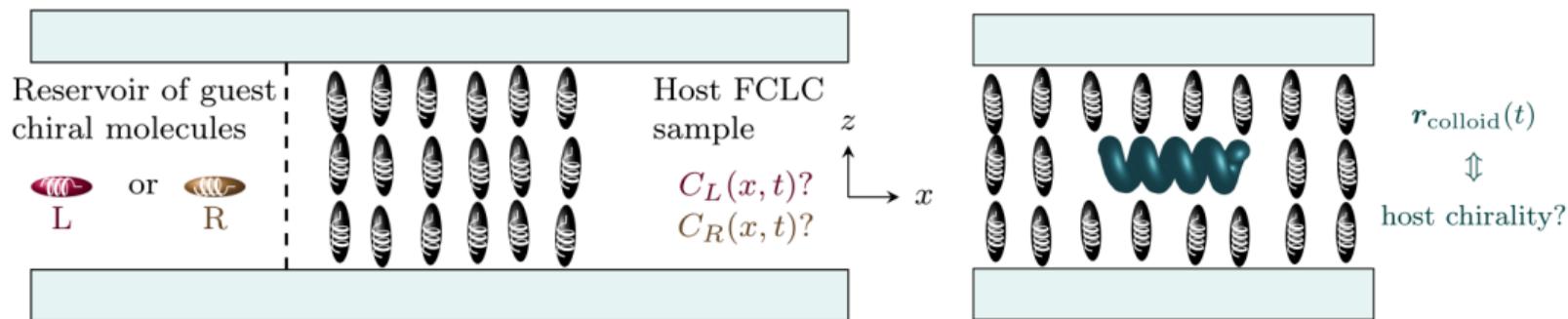
## Towards topological optomechanics



## Towards topological optomechanics



## CNRS{5}: Diffusive transport properties in chiral guest-host systems



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## Realization and Application of Topological Defect Patterns in Soft and Living Matter

### Guest Editors

Dr. Simon Čopar, Dr. Guilhem Poy, Prof. Dr. Anupam Sengupta

### Deadline

01 June 2021

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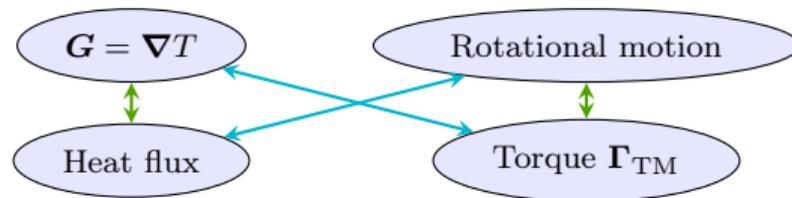
# Special Issue

Invitation to submit

Thank you for your attention!

## Leslie interpretation of the Lehmann experiment

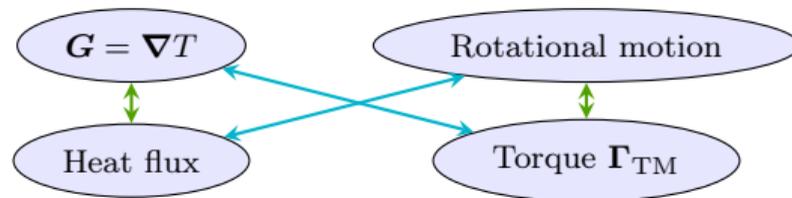
First explanation by Leslie in 1968:



- Existence, in a cholesteric phase, of a torque on the director:  
 $\Gamma_{TM} = \nu \mathbf{n} \times [\mathbf{n} \times \mathbf{G}]$ , with  $\nu$  the Leslie thermomechanical coefficient.

## Leslie interpretation of the Lehmann experiment

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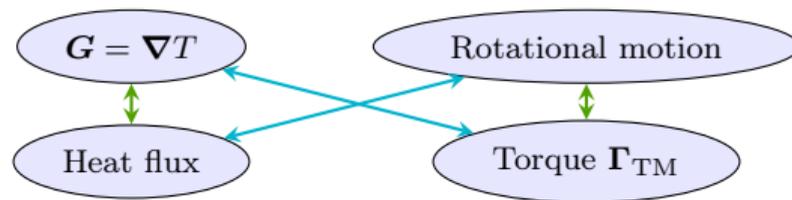


- Existence, in a cholesteric phase, of a torque on the director:  
 $\Gamma_{TM} = \nu \mathbf{n} \times [\mathbf{n} \times \mathbf{G}]$ , with  $\nu$  the Leslie thermomechanical coefficient.
- As in a wind turbine, essential role of the chirality:  
 no rotation predicted in a nematic phase.

F. M. Leslie. *Proceedings of the Royal Society A*, 307:359–372, 1968

## Leslie interpretation of the Lehmann experiment

First explanation by Leslie in 1968:

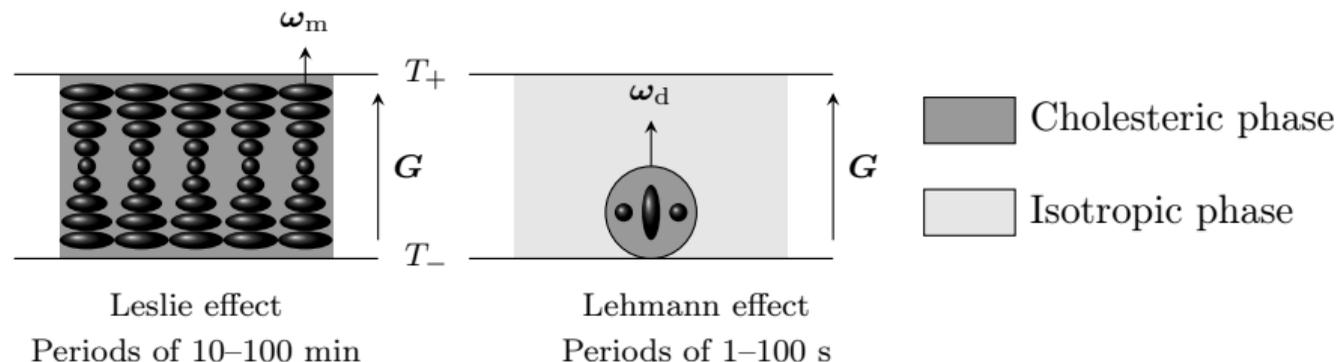


## Leslie paradigm

The rotation of the texture in the Lehmann experiment is due to the Leslie thermomechanical torque  $\Gamma_{TM}$

## Lehmann vs. Leslie experiment

Oswald &amp; Dequidt, 2008-2014:

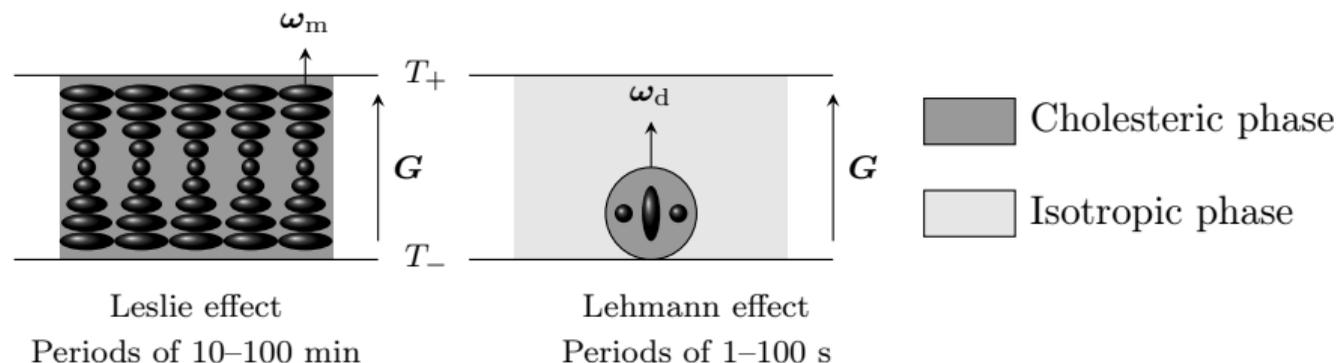


P. Oswald and A. Dequidt. *Physical Review Letters*, 100:217802, 2008

P. Oswald. *Europhysics Letters*, 108:36001, 2014

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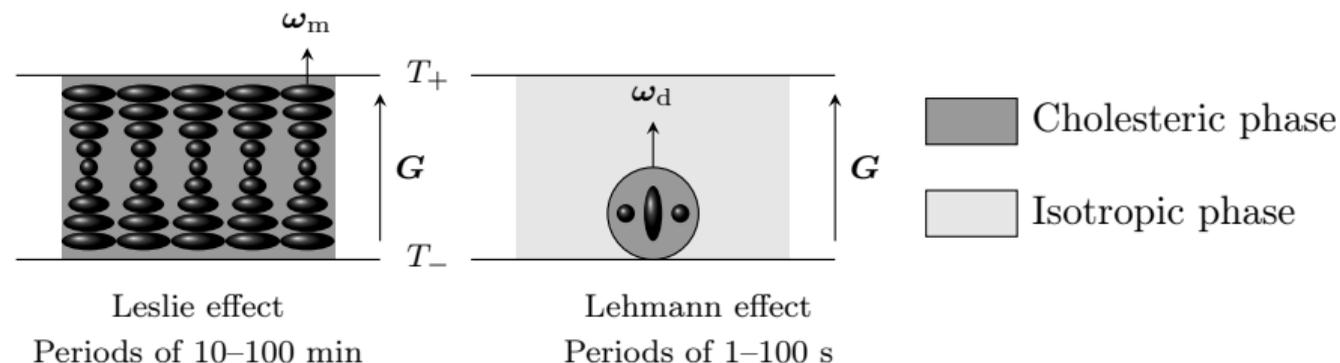
$\omega_d$  and  $\omega_m$  sometimes of opposite signs!

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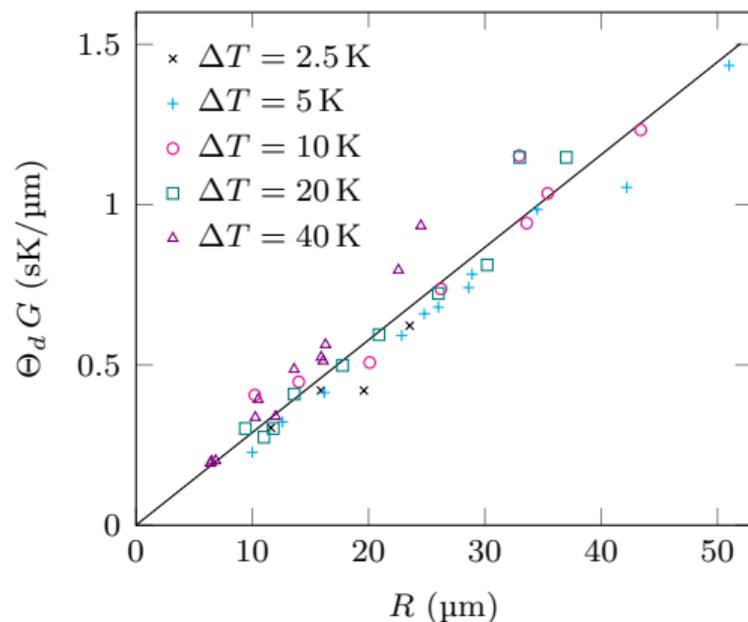
$\omega_d$  and  $\omega_m$  sometimes of opposite signs!

Leslie effect  $\neq$  Lehmann effect?

P. Oswald and A. Dequidt. *Physical Review Letters*, 100:217802, 2008

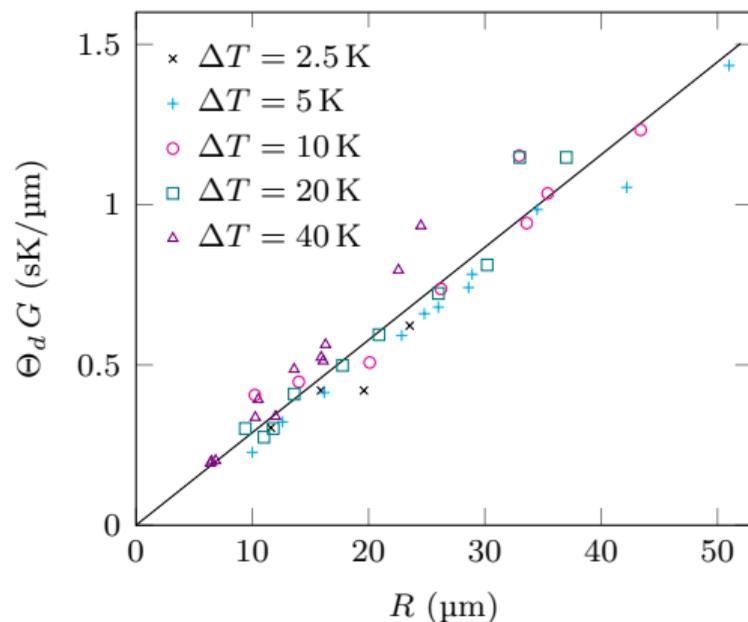
P. Oswald. *Europhysics Letters*, 108:36001, 2014

## Rotation periods of SSY droplets



- Angular velocity  $\omega_d = 2\pi/\Theta_d$  proportional to  $G$ .

## Rotation periods of SSY droplets



- Angular velocity  $\omega_d = 2\pi/\Theta_d$  proportional to  $G$ .
- Period  $\Theta_d$  proportional to the radius  $R$ .